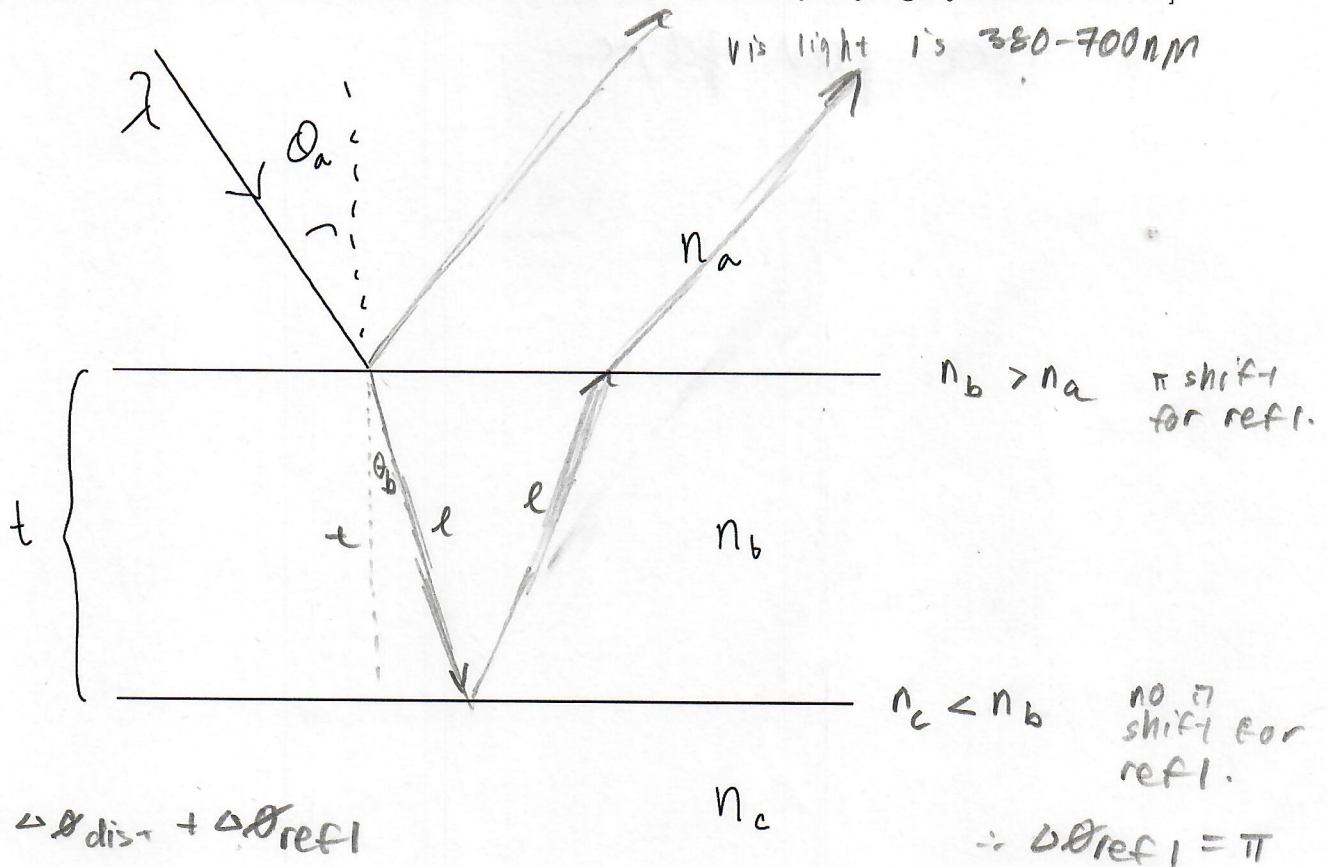


Problem 1

20 points

Consider a thin film of thickness $t = 2.30 \times 10^{-6}$ m and index of refraction $n_b = 1.20$. The film is resting on a material of index of refraction $n_c = 1.13$, and its top face is exposed to air $n_a = 1.00$. What is the shortest wavelength of visible light that will interfere destructively when incident on the film at angle $\theta_a = 22.0^\circ$ from the normal? Give your answer in nanometers, to three significant figures. [Note: you may ignore the fact that the wavelength will change upon refraction; this will only very slightly affect the answer.]



$$\Delta\theta_{tot} = \Delta\theta_{dist} + \Delta\theta_{refl}$$

n_c

$$\therefore \Delta\theta_{refl} = \pi$$

$$\frac{2\pi}{\lambda} \Delta x + \pi = (2m+1)\pi$$

$$\frac{4\pi t}{\lambda} \cos\theta_b + \pi = (2m+1)\pi$$

$$\frac{4\pi t}{\lambda} \cos\theta_b = 2m$$

$$\frac{2t}{\lambda} \cos\theta_b = m$$

$$\frac{2t}{\lambda \cos\theta_b} = m$$

$$\lambda = \frac{2t}{m \cos\theta_b} \text{ as } m \uparrow, \lambda \downarrow$$

$$\text{try } \lambda = 380 \times 10^{-9} \text{ m}$$

$$m = \frac{2t}{\lambda \cos\theta_b} = 12.74$$

13 would be too short of a wavelength so $m = 12$

find Δx

$$n_a \sin\theta_a = n_b \sin\theta_b$$

$$\Delta x = 2t$$

$$\cos\theta_b = \frac{t}{\ell}$$

$$\begin{aligned} \theta_b &= \sin^{-1}\left(\frac{\sin\theta_a}{n_b}\right) \\ &= \sin^{-1}\left(\frac{\sin(22)}{1.2}\right) \\ &= 18.19 \end{aligned}$$

$$\begin{aligned} \ell &= \frac{t}{\cos\theta_b} \\ &= \frac{2.3 \times 10^{-6}}{\cos(18.19)} \\ &= 2.421 \times 10^{-6} \end{aligned}$$

464nm

$$\lambda_{smallest} = \frac{2\ell}{12} = \frac{\ell}{6} = 4.03497 \times 10^{-7} \text{ m}$$

403.5 nm

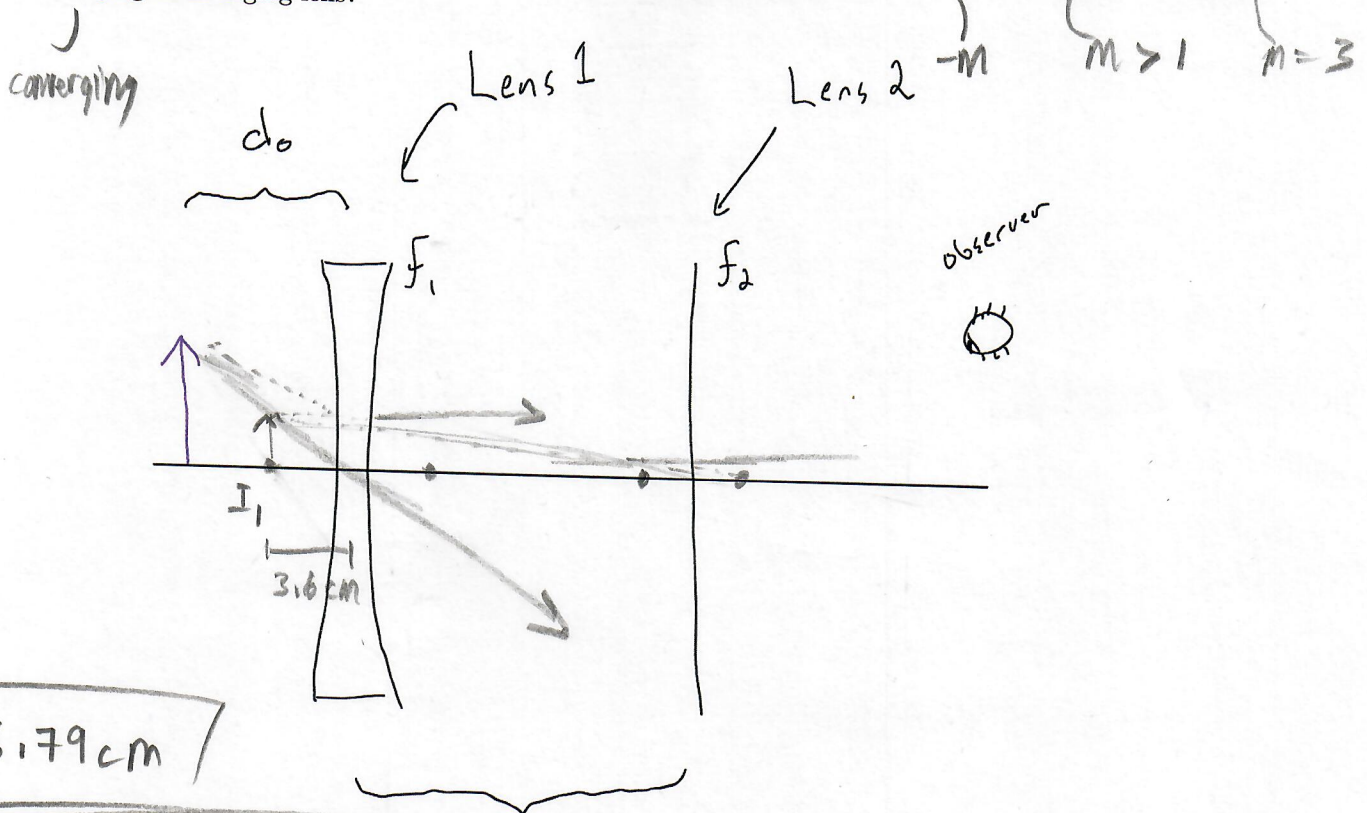
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Problem 2

20 points

Consider the thin lens (Lens 1) of focal length $f_1 = -2.00$ cm shown below. An object is placed a distance $d_o = 4.50$ cm to the left of the lens. A second lens (Lens 2, drawn as a single vertical line in the diagram) is then placed $D = 5.00$ cm to the right of Lens 1. What is the focal length Lens 2 must have such that the final image of the object seen through the two lenses is inverted and enlarged by a factor of three with respect to the object? Give your answer in centimeters, to three significant figures. Is Lens 2 a converging or diverging lens?



$f_2 = 6.79 \text{ cm}$

converging lens, $f > 0$

$$\frac{1}{f} = \frac{1}{s} + \frac{1}{s'}$$

$$\frac{s-f}{sf} = \frac{1}{s'}$$

$$s' = \frac{sf}{s-f}$$

$$= \frac{4.5(2)}{4.5-2}$$

$$= -3.6$$

$$m = -\frac{s'}{s}$$

$$= -\frac{3.6}{4.5}$$

$$= 0.8$$

$$m = \frac{y_1'}{y_1}$$

$$s_1 = D - s'$$

$$= 5 + 3.6$$

$$y_1' = 0.8 y_1$$

$$0.8 m_{12} = \frac{y_2}{y_1} = -3$$

$$m_{12} = -3.75$$

for: f_2 known: s_1, m_{10}, m_{20}

$$\frac{y_2}{y_1} = -3 = m_{12} = -\frac{s_2}{s_1}$$

$$s_2 = -m_{12} s_1$$

$$= 3.75(8.1)$$

$$= 30.375$$

$$f_2 = \frac{30.375(8.1)}{8.1+30.375}$$

$$= 6.789 \text{ cm}$$

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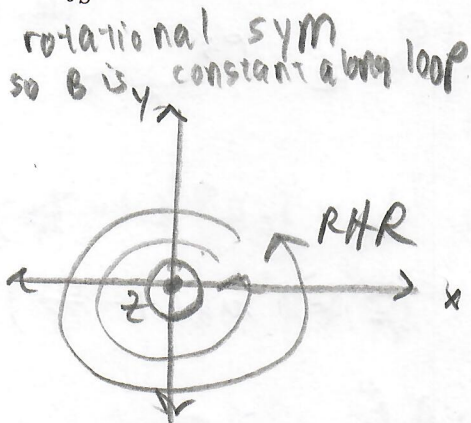
Problem 3

20 points

Consider an infinitely long cylindrical wire of radius R centered on the z axis. The current density in the wire varies with the distance ρ from the axis as $\vec{J}(\rho) = J_0 \left(1 - \frac{\rho}{R}\right) \hat{z}$, where J_0 is a constant.

(a): 10 points

Find the magnetic field everywhere in space. Recall that current through a surface S is defined as $I = \int_S \vec{J} \cdot d\vec{A}$.



$$d\vec{A} = 2\pi\rho d\rho$$

$$I_{\text{enc}} = \int_0^{\rho} J_0 \left(1 - \frac{\rho'}{R}\right) \cdot 2\pi\rho' d\rho'$$

(inside wire)

$$= 2\pi J_0 \left[\frac{1}{2}\rho'^2 - \frac{1}{3R}\rho'^3 \right]_0^{\rho}$$

$$= 2\pi J_0 \left[\frac{1}{2}\rho^2 - \frac{1}{3R}\rho^3 \right]$$

$$\oint \vec{B} \cdot d\vec{l} = \mu_0 I_{\text{enc}}$$

amp. curve of radius r

$$I_{\text{enc}} = \int_0^R J_0 \left(1 - \frac{\rho}{R}\right) \cdot 2\pi\rho d\rho$$

(total)

$$= J_0 2\pi \int_0^R \rho - \frac{1}{R}\rho^2 d\rho$$

$$= 2\pi J_0 \left[\frac{1}{2}\rho^2 - \frac{1}{3R}\rho^3 \right]_0^R$$

$$= 2\pi J_0 \left[\frac{1}{2}R^2 - \frac{1}{3R}R^3 \right]$$

$$= 2\pi J_0 \left[R^2 \left(\frac{1}{2} - \frac{1}{3} \right) \right]$$

$$= 2\pi J_0 R^2 \cdot \frac{1}{6}$$

$$= \frac{1}{3}\pi J_0 R^2$$

$$\oint \vec{B} \cdot d\vec{l} = B \cdot 2\pi\rho$$

inside

$$B \cdot 2\pi\rho = 2\pi J_0 \left[\frac{1}{2}\rho^2 - \frac{1}{3R}\rho^3 \right] \mu_0$$

$$B = \frac{J_0 \mu_0}{\rho} \left[\frac{1}{2}\rho^2 - \frac{1}{3R}\rho^3 \right]$$

$$B = \mu_0 J_0 \left[\frac{1}{2}\rho - \frac{1}{3R}\rho^2 \right]$$

outside

$$B \cdot 2\pi\rho = \frac{1}{3}\pi J_0 R^2 \mu_0$$

$$B = \frac{1}{6\rho} J_0 R^2 \mu_0$$

direction: rotate \hat{z} around z -axis as shown in diagram

(b): 10 points

Suppose now that the current density falls off as a function of time:

$$\vec{J}(\rho, t) = J_0 \left(1 - \frac{\rho}{R}\right) \left(1 - \frac{t}{\tau}\right) \hat{z}$$

where τ is a constant and $0 \leq t \leq \tau$. Calculate the electric field everywhere outside the wire, in terms of μ_0 , ϵ_0 , J_0 , R , τ , and any spatial coordinates and/or time. [Hint: the electric field will point along the z axis. Use Faraday's law and integrate around a rectangle whose base and top are parallel to the z axis.] [Note: you will only be able to calculate the electric field up to an arbitrary constant; this is ok.]

$$\oint \vec{E} \cdot d\vec{\ell} = - \frac{d\Phi_E}{dt}$$

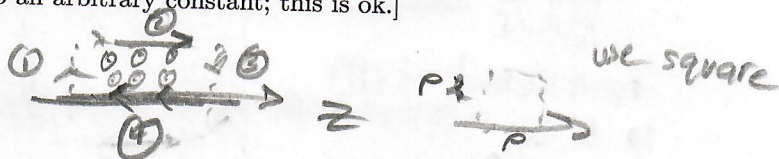
1) find B

$$\begin{aligned} I_{\text{enc}} &= \int_0^R J_0 \left(1 - \frac{\rho}{R}\right) \left(1 - \frac{t}{\tau}\right) 2\pi\rho d\rho \\ &= \left(1 - \frac{t}{\tau}\right) \int_0^R J_0 \left(1 - \frac{\rho}{R}\right) 2\pi\rho d\rho \\ &= \frac{1}{3} \pi J_0 R^2 \left(1 - \frac{t}{\tau}\right) \end{aligned}$$

2) find $\Phi_B(t)$

$$\begin{aligned} d\Phi_B &= \vec{B} \cdot d\vec{A} \\ \Phi_B &= \int \vec{B} \cdot d\vec{A} \\ &= \int \vec{B} \cdot \rho d\vec{B} \\ &= \int \frac{1}{6} J_0 R^2 \mu_0 \left(1 - \frac{t}{\tau}\right) dA \\ &= \frac{1}{6} J_0 R^2 \mu_0 \rho \left(1 - \frac{t}{\tau}\right) \end{aligned}$$

$$\frac{d\Phi_B}{dt} = - \frac{1}{6\tau} J_0 R^2 \mu_0 \rho$$



$$\begin{aligned} B \cdot 2\pi\rho &= \frac{1}{3} \pi J_0 R^2 \left(1 - \frac{t}{\tau}\right) \mu_0 \\ B &= \frac{1}{6\rho} J_0 R^2 \mu_0 \left(1 - \frac{t}{\tau}\right) \end{aligned}$$

3) find $\oint \vec{E} \cdot d\vec{\ell}$

\vec{E} is in \hat{z} direc so
 $\oint = 0$ along sides 1, 3
 E is constant along the straight line. \odot along side 4 b/c no mag field there (no charge encl @ middle of wire)
 $\therefore \oint \vec{E} \cdot d\vec{\ell} = E\rho$

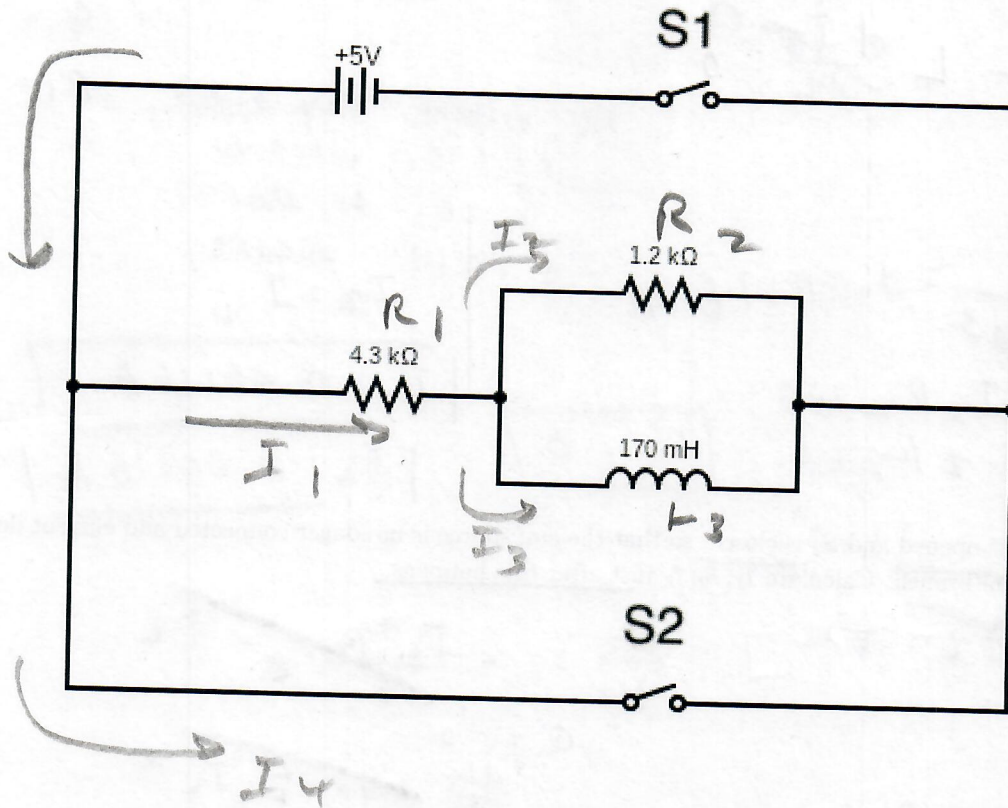
$$E\rho = \frac{1}{6\tau} J_0 R^2 \mu_0 \rho$$

$$\boxed{E = \frac{1}{6\tau} J_0 \mu_0 R^2}$$

Problem 4

20 points

Consider the circuit drawn below. Let I_1 be the current flowing through the $4.3\text{ k}\Omega$ resistor, I_2 be the current flowing through the $1.2\text{ k}\Omega$ resistor, and I_3 be the current flowing through the 170 mH inductor.



(a): 5 points

Suppose S_1 has been closed for a long time, ($-\infty < t < 0$), and the emf source drives a steady current. Calculate I_1, I_2, I_3 .

$$\downarrow \frac{dI}{dt} = 0$$

$$\mathcal{E} - I_1 R_1 - I_2 R_2 = 0$$

$$I_1 = I_2 + I_3$$

$$\mathcal{E} - I_1 R_1 - L \frac{dI_3}{dt} = 0$$

$$\mathcal{E} = 5$$

$$R_1 = 4.3 \text{ k}\Omega$$

no voltage drop across inductor when $\frac{dI}{dt} = 0$ so short circuit occurs

$$\mathcal{E} = I_1 R_1$$

$$I_1 = \frac{5}{4.3 \times 10^3} = 0.001163 \text{ A}$$

$$I_3 = I_1$$

$$\mathcal{E} - \mathcal{E} - I_2 R_2 = 0$$

$$I_2 R_2 = 0$$

$$I_2 = 0$$

$$I_1 = 0.00116 \text{ A}$$

$$I_3 = 0.00116 \text{ A}$$

(b): 5 points

At time $t = 0$, S_1 is opened and S_2 is closed, so that the emf source is no longer connected and current flows through the bottom branch. Calculate I_1, I_2, I_3 just after this happens.

~~$$I_3 = I_3^* e^{-\frac{R}{L}t} \quad \frac{dI_3}{dt} = -\frac{R}{L} I_3^* e^{-\frac{R}{L}t}$$~~

from part a

$$I_3 = I_3^* e^{0} = I_3^*$$

@ $t = 0$

~~$$\frac{dI_3}{dt} = -\frac{R}{L} I_3^*$$~~

$$I_1 = I_2 + I_3$$

$$L \frac{dI_3}{dt} = I_1 R_1$$

~~$$\int \frac{1}{I_1} dI_3 \int \frac{R_1}{L} dt \quad ???$$~~

$$L \frac{dI_3}{dt} - I_1 R_1 = 0$$

~~$$I_1 R_1 = L \frac{dI_3}{dt} \quad \int (I_1) = I_2 + I_3$$~~

~~$$I_1 = -\frac{L}{R_1} \cdot \frac{R}{L} I_3^* = I_2 + I_3$$~~

current is dec so this term is +

@ $t = 0$, no time to discharge so current does not change?

$$I_1 = 0.00116 \text{ A} \quad I_3 = 0.00116 \text{ A} \quad I_2 = 0$$

(c): 5 points

Calculate I_1, I_2, I_3 in the limit $t \rightarrow \infty$.

as $t \rightarrow \infty$, $\frac{dI_3}{dt} \rightarrow 0$

system is discharged
so $I_3 = 0$

$$L \frac{dI_3}{dt} - I_1 R_1 = 0 \quad I_1 = 0$$

$$I_1 = I_2 + I_3$$

$$\boxed{I_1 = 0 \text{ A}, I_2 = 0 \text{ A}, I_3 = 0 \text{ A}}$$

(d): 5 points

Calculate the total energy dissipated in all resistors during the time period $0 < t < \infty$.

$$U = \frac{1}{2} L i^2 \quad \rightarrow \text{must be equal to all energy present @ the start}$$

$$\begin{aligned} \text{energy stored in inductor} &= \frac{1}{2} (170 \times 10^{-3}) (0.00116)^2 \\ &= \boxed{1.14376 \times 10^{-7} \text{ J}} \end{aligned}$$

Problem 5

10 points

(a): 5 points

Consider the following vector field:

$$\vec{V}(r, \theta, \phi) = \alpha \frac{1}{r^2} \hat{r},$$

where α is a constant with units $\text{T}\cdot\text{m}^2$, r is the distance from the origin, and \hat{r} is a unit vector pointing away from the origin. According to Maxwell's equations, could this represent a magnetic field? Explain.

$$\nabla \times \vec{B} = \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t} \quad \nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} \quad \nabla \cdot \vec{E} = \frac{\rho}{\epsilon_0} \quad \nabla \cdot \vec{B} = 0$$

$$\nabla \vec{V} = \left\langle \frac{\partial \vec{V}}{\partial r}, \frac{\partial \vec{V}}{\partial \theta}, \frac{\partial \vec{V}}{\partial \phi} \right\rangle$$

$$= \left\langle -\frac{2\alpha}{r^3}, 0, 0 \right\rangle$$

$$\nabla \vec{V} \cdot \vec{V} = -\frac{2\alpha}{r^3} \left(\alpha \frac{1}{r^2} \right) = -\frac{2\alpha^2}{r^5} \neq 0$$

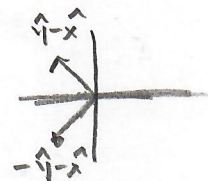
this vector field does not meet the divergence condition and \therefore cannot be a magnetic field

(b): 5 points

Consider the following electric and magnetic fields:

$$\vec{E} = \frac{E_0}{\sqrt{2}} \cos(\omega(t - z/c)) (\hat{y} - \hat{x})$$

$$\vec{B} = \frac{E_0}{c\sqrt{2}} \cos(\omega(t - z/c)) (-\hat{y} - \hat{x})$$



Can these fields constitute a traveling electromagnetic plane wave in vacuum? If yes, prove it. If not, explain why not.

• $B_0 = \frac{E_0}{c}$ - amplitudes are proportional

• $\hat{y} - \hat{x}$ is \perp to $-\hat{y} - \hat{x}$ and each is \perp to the z prop. direction

• waves travel in the same direction (z) and are in phase (same cos arg.)

the two fields meet all necessary conditions so they can constitute a traveling EM plane wave

Problem 6

10 points

Consider an LRC series circuit driven with an ac source $v(t) = V_0 \cos(\omega t)$. You may use without proof the impedance and phase of an LRC series circuit.

(a): 5 points

Calculate the total energy dissipated in the resistor over one cycle if the system is driven at its resonant frequency.

$$X_L = X_C$$

$$\omega_0 = \frac{1}{\sqrt{LC}}$$

$$P_R = i^2 R$$

$$P_{av} = \frac{1}{2} I^2 R$$

$$\tan \theta = \frac{X_L - X_C}{R}$$

$$\theta = 0, \quad X_L - X_C = 0$$

$$\therefore i(t) = I \cos(\omega_0 t)$$

$$\cos^2 x + \sin^2 x = 1$$

$$\cos 2x = \cos^2 x - \sin^2 x$$

$$= \cos^2 x - (1 - \cos^2 x)$$

$$\cos 2x = 2\cos^2 x - 1$$

$$\frac{1}{2} (\cos 2x + 1) = \cos^2 x$$

$$V_0 = I \sqrt{R^2 + (X_L - X_C)^2} = I Z = I R$$

$$i(t) = I \cos(\omega_0 t)$$

$$P_R = I^2 R \cos^2(\omega_0 t)$$

$$T = \frac{2\pi}{\omega}$$

$$P_{tot} = \int_0^{\frac{2\pi}{\omega_0}} I^2 R \cos^2(\omega_0 t) dt$$

$$= I^2 R \int_0^{\frac{2\pi}{\omega_0}} \cos^2(\omega_0 t) dt$$

$$= I^2 R \int_0^{\frac{2\pi}{\omega_0}} \frac{1}{2} (\cos(2\omega_0 t) + 1) dt$$

$$= \frac{1}{2} I^2 R \int_0^{\frac{2\pi}{\omega_0}} \cos(2\omega_0 t) + 1 dt$$

$$= \frac{1}{2} I^2 R \left[\frac{1}{2\omega_0} \sin(2\omega_0 t) + t \right]_0^{\frac{2\pi}{\omega_0}}$$

$$= \frac{1}{2} I^2 R \left[\frac{1}{2\omega_0} \sin(4\pi) + \frac{2\pi}{\omega_0} - 0 \right]$$

$$= \frac{\pi}{\omega_0} I^2 R \quad [\text{just } P_{av} \cdot T]$$

$$= \frac{R\pi}{\omega_0} \left(\frac{V_0}{R} \right)^2$$

$$= \frac{V_0^2 \pi}{R \omega_0}$$

$$P_{tot} = \frac{V_0^2 \pi}{R \omega_0}$$

(b): 5 points

Calculate the total energy dissipated in the resistor over one cycle if the system is driven at twice its resonant frequency. Is it greater than or less than the energy found in part (a)?

$$Z_{LRC} = \sqrt{R^2 + (L\omega - (1/\omega C))^2} \quad \omega = 2\omega_0$$

$$2\omega_0 L - \frac{1}{2\omega_0 C}$$

$$\omega_0 = \frac{1}{\sqrt{LC}}$$

$$= \frac{2L}{\sqrt{LC}} - \frac{\sqrt{LC}}{2C}$$

$$I Z_{LRC} = V_0$$

$$I = \frac{V_0}{Z_{LRC}}$$

$$= \frac{4LC}{2C\sqrt{LC}} - \frac{LC}{2C\sqrt{LC}}$$

$$P_{av} = \frac{1}{2} I^2 R$$

$$Z_{LRC} = \sqrt{\frac{3LC}{2C\sqrt{LC}}}$$

$$= \frac{1}{2} \left(\frac{V_0}{\sqrt{R^2 + \frac{9L}{4C}}} \right)^2 R$$

$$= \frac{3L}{2\sqrt{LC}}$$

$$= \frac{1}{2} R \cdot \frac{V_0^2}{R^2 + \frac{9L}{4C}}$$

$$Z_{LRC} = \sqrt{R^2 + \left(\frac{3L}{2\sqrt{LC}} \right)^2}$$

$$P_{tot} = P_{av} \cdot \frac{2\pi}{2\omega_0}$$

$$= \sqrt{R^2 + \frac{9L^2}{4LC}}$$

$$= P_{av} \cdot \frac{\pi}{\omega_0}$$

$$= \sqrt{R^2 + \frac{9L}{4C}}$$

$$= \frac{R V_0^2}{2(R^2 + \frac{9L}{4C})} \cdot \frac{\pi}{\omega_0}$$

Since R is greater than R
 $\frac{2(R^2 + \frac{9L}{4C})}$ more energy is dissipated in resistor when system is @ resonant freq

$$= \frac{V_0^2 \pi}{\omega_0} \cdot \frac{R}{2(R^2 + \frac{9L}{4C})}$$

conclusion makes sense
 as ω_0 minimizes impedance
 L and C contrib. from

$2R^2 + \frac{9L}{2C} > R^2$
 b/c $\frac{9L}{2C}$ is always positive. also $\times 2 R^2$

compare to $\frac{V_0^2 \pi}{\omega_0} \cdot \frac{1}{R}$

$$\frac{1}{R} > \frac{R}{2(R^2 + \frac{9L}{4C})}$$

$$\frac{1}{R^2} > \frac{1}{2(R^2 + \frac{9L}{4C})}$$