20 points

(equation # in textbook)

Consider a thin film of thickness $t = 2.30 \times 10^{-6}$ m and index of refraction $n_b = 1.20$. The film is resting on a material of index of refraction $n_c = 1.13$, and its top face is exposed to air $n_a = 1.00$. What is the shortest wavelength of visible light that will interfere destructively when incident on the film at angle $\theta_a = 22.0^{\circ}$ from the normal? Give your answer in nanometers, to three significant figures. [Note: you may ignore the fact that the wavelength will change upon refraction; this will only very slightly affect the answer.]



wave 1:

-shifts by a half cycle because nb> na

wave 2:

- does not shift because ne < nb

so for destructive interference,

 $2x = m\lambda$ (m = D, 1, 2, ...) (35.186)

lignore that & changes upon refraction

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2+ = mλ WS OL 入= 2+ m (OS Ob Snell's Law: Nasinda = nisinde (33.4) $\Theta_{L} = \sin^{-1}\left(\frac{Na\sin\Theta_{A}}{NL}\right)$ $= \sin^{-1} \left(\frac{1.00(\sin 22)}{1.20} \right)$ = 18.19018° $\lambda = \frac{2(2300)}{m \cos(18.19018)}$ $= 4841.975 \quad (m = 1, 2, 3, ...)$

smallest wavelength of visible light is 380 nm

if
$$m = 13$$
,
 $\lambda = 372 \text{ nm} (\text{not visible})$
if $m = 12$,
 $\lambda = 403 \text{ nm}$

20 points

(equation # in textbook)

Consider the thin lens (Lens 1) of focal length $f_1 = -2.00$ cm shown below. An object is placed a distance $d_o = 4.50$ cm to the left of the lens. A second lens (Lens 2, drawn as a single vertical line in the diagram) is then placed D = 5.00 cm to the right of Lens 1. What is the focal length Lens 2 must have such that the *final* image of the object seen through the two lenses is inverted and enlarged by a factor of three with respect to the object? Give your answer in centimeters, to three significant figures. Is Lens 2 a converging or diverging lens?



be a

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lend 2:

$$\frac{1}{138441+500} + \frac{1}{5_{2}} = \frac{1}{f_{2}}$$

$$\frac{1}{638441} + \frac{1}{5_{2}} = \frac{1}{f_{2}}$$

$$m = m_{1}m_{2}$$

$$m_{2} = \frac{m}{m_{1}}$$

$$= -\frac{3}{0.30769}$$

$$= -9.75$$

$$m_{2} = -5_{2}'$$

$$s_{2}' = -m_{2}S_{2}$$

$$= -(-9.75)(6.38461)$$

$$= 62.25$$

$$f_{2} = \frac{1}{\frac{1}{6.38461} + \frac{1}{62.25}}$$

$$= 5.79 \text{ cm}}$$
diverging tenses have negative focal lengths, and f_{2} is positive, so tens 2 must be a
Converging tenses

20 points

(equation # in textbook)

Consider an infinitely long cylindrical wire of radius R centered on the z axis. The current density in the wire varies with the distance ρ from the axis as $\vec{J}(\rho) = J_0 \left(1 - \frac{\rho}{R}\right) \hat{z}$, where J_0 is a constant.

(a): 10 points

Find the magnetic field everywhere in space. Recall that current through a surface S is defined as $I = \int_{S} \vec{J} \cdot d\vec{A}$.

because of symmetry, we can use Ampere's Law $\oint_{3} \vec{B} \cdot d\vec{1} = M_0 \text{Ience}(29.19)$ $B \oint_{3} s d\vec{1} = M_0 \text{Ience}$ $B 2 \pi P = M_0 \text{Ience}$ $B = \frac{M_0 \text{Ience}}{2 \pi P}$

$$\begin{split} & \text{if } \rho < R, \\ & \text{I} = \int_{0}^{\rho} J_{0} \left(1 - \frac{\rho}{R} \right) 2 \pi \rho \, d\rho \\ & = 2\pi J_{0} \int_{0}^{\rho} \rho - \frac{\rho^{2}}{R} \, d\rho \\ & = 2\pi J_{0} \left(\frac{\rho^{2}}{2} - \frac{\rho^{3}}{3R} \right) \\ & = 2\pi J_{0} \left(\frac{\rho^{2}}{2} - \frac{\rho^{3}}{3R} \right) \\ & = 2\pi J_{0} \rho^{2} \left(\frac{1}{2} - \frac{\rho}{3R} \right) \\ & = 2\pi J_{0} \rho^{2} \left(\frac{1}{2} - \frac{\rho}{3R} \right) \\ & = \frac{\pi J_{0} R^{2} \left(\frac{1}{2} \right)_{3}}{2\pi r} \\ & = \frac{\pi J_{0} R^{2} \left(\frac{1}{2} \right)_{3}}{3(2\pi r)} \\ & = \frac{\pi J_{0} \sigma J_{0} \rho^{2} \left(\frac{1}{2} - \frac{\rho}{3R} \right) \hat{x} \\ & = \frac{\pi J_{0} \sigma J_{0} \rho^{2} \left(\frac{1}{2} - \frac{\rho}{3R} \right) \hat{x} \\ & = \frac{\pi J_{0} \sigma J_{0} \rho^{2} \hat{x}}{b \rho} \end{split}$$

Ζ

÷Χ

(b): 10 points

Suppose now that the current density falls off as a function of time:

$$\vec{J}(\rho,t) = J_0 \left(1 - \frac{\rho}{R}\right) \left(1 - \frac{t}{\tau}\right) \hat{z} \qquad \checkmark \boldsymbol{P} > \boldsymbol{R}$$

where τ is a constant and $0 \le t \le \tau$. Calculate the electric field everywhere outside the wire, in terms of $\mu_0, \epsilon_0, J_0, R, \tau$, and any spatial coordinates and/or time. [Hint: the electric field will point along the z axis.] Use Faraday's law and integrate around a rectangle whose base and top are parallel to the z axis.] [Note: you will only be able to calculate the electric field up to an arbitrary constant; this is ok.]

$$\begin{split} \oint \mathbf{E} \cdot d\mathbf{\overline{I}} &= -\frac{d\Phi_{B}}{dt} (29.10) \\ P > R, S0 \\ B &= \underline{M \cup J \circ R^{2}} \left(1 - \frac{1}{T} \right) (\text{from part } n) \\ \overline{\Phi}_{B} &= \int B \cdot dA (27.6), \ dA &= L \ dP \\ &= \underline{M \circ J \circ R^{2}} \left(1 - \frac{1}{T} \right) \int_{P}^{P+W} \frac{1}{P} \ dP \\ &= \underline{M \circ J \circ R^{2}} \left(1 - \frac{1}{T} \right) \ln \left(\frac{P+W}{P} \right) \\ e^{S} &= \frac{M \circ J \circ R^{2}}{6T} \ln \left(\frac{P+W}{P} \right) \\ e^{S} &= \frac{M \circ J \circ R^{2}}{6T} \ln \left(\frac{P+W}{P} \right) \\ e^{S} &= \frac{M \circ J \circ R^{2}}{6T} \ln \left(\frac{P+W}{P} \right) \hat{z} \\ &= \frac{M \circ J \circ R^{2}}{12LT} \ln \left(\frac{P+W}{P} \right) \hat{z} \end{split}$$

20 points

(equation # in textbook)

Consider the circuit drawn below. Let I_1 be the current flowing through the 4.3 k Ω resistor, I_2 be the current flowing through the 1.2 k Ω resistor, and I_3 be the current flowing through the 170 mH inductor.



Vsing Kirchhoff's Rules,
1)
$$I_1 = I_2 + I_3$$

2) $E - I_1 R_1 - I_2 R_2 = 0$
3) $E - I_1 R_1 - L d I_3 = 0$ (30.13)
 $d + d = 0$

Rombes

(a): 5 points

Suppose S_1 has been closed for a long time, $(-\infty < t < 0)$, and the emf source drives a steady current. Calculate I_1, I_2, I_3 .



(b): 5 points

At time t = 0, S_1 is opened and S_2 is closed, so that the emf source is no longer connected and current flows through the bottom branch. Calculate I_1, I_2, I_3 just after this happens.

now using kirchhoff's kiles,
3)
$$-I_1R_1 - L \frac{dI_3}{dt} = 0$$

 $T = \frac{dI_3}{dt} = -I_1R_1 - Rea$
 $\frac{dI_3}{at} = -I_1R_1 - Rea$
 $\frac{dI}{at} = -\frac{IRea}{L}$
 $\int \frac{dI}{dI} = -\int \frac{Rea}{L} dt$
 $I_1 = 1.16 \times 10^{-3} A$
 $I_2 = D$
 $\int \frac{dI}{I} = -\int \frac{Rea}{L} dt$
 $I_3 = 1.16 \times 10^{-3} A$
 $In(I) = -\frac{Rea}{L} t$
 $I = I_0 e^{-Re}/L t$ (30.18)

(c): 5 points

Calculate I_1, I_2, I_3 in the limit $t \to \infty$.

$$I = I_0 e^{-Ry} L^{\dagger}$$

$$\lim_{t \to \infty} I_0 e^{-Ry} L^{\dagger}$$

$$= I_0 e^{-\infty}$$

$$= I_0 (0)$$

$$= 0$$

(d): 5 points

Calculate the total energy dissipated in all resistors during the time period $0 < t < \infty$.

evergy dissipated through both resistors
energy dissipated =
$$I^{2}Req$$
 (30.19)
= $Ren \int_{0}^{\infty} (I_{0} e^{-Req/Lt})^{2} dt$
= $Ren I_{0}^{2} \int_{0}^{\infty} (e^{-Req/Lt})^{2} dt$
= $Ren I_{0}^{2} \int_{0}^{\infty} (e^{-Req/Lt})^{2} dt$
= $\frac{I_{0}^{2} L}{2Rent} = \frac{I_{0}^{2} L}{2Rent} (0-1)$
= $\frac{I_{0}^{2} L}{2}$
= $(1.16 \times 10^{-3})^{2} (0.17)$
= $1.14 \times 10^{-7} J$

10 points

(a): 5 points

(equation # in textbook)

Consider the following vector field:

$$\vec{V}(r,\theta,\phi) = \alpha \frac{1}{r^2}\hat{r},$$

where α is a constant with units T·m², r is the distance from the origin, and \hat{r} is a unit vector pointing away from the origin. According to Maxwell's equations, could this represent a magnetic field? Explain.

the only Maxwell equation that restricts whether fields
can be magnetic fields is
$$\nabla \cdot \vec{B} = 0$$
$$= \frac{1}{r^2} \frac{d}{dr} \left(p \neq \alpha \pm \right) \neq divergence in spherical coordinates$$
$$= 0 \checkmark$$

Yes, it satisfies the Maxwell equation

(b): 5 points

Consider the following electric and magnetic fields:

$$\vec{E} = \frac{E_0}{\sqrt{2}} \cos \left(\omega(t - z/c) \right) (\hat{y} - \hat{x})$$

$$\vec{B} = \frac{E_0}{c\sqrt{2}} \cos \left(\omega(t - z/c) \right) (-\hat{y} - \hat{x}).$$
direction
direction

- < -1, 1, 0 > is in this

Can these fields constitute a traveling electromagnetic plane wave in vacuum? If yes, prove it. If not, explain why not.

Conditions
1.
$$\vec{E}$$
 and \vec{B} are \perp (from $\oint \vec{E} \cdot dA = \oint \vec{B} \cdot dA = 0$)
 $\langle -1, 1, 0 \rangle \cdot \langle -1, -1, 0 \rangle = 1 - 1 = 0$, so they are $\perp \checkmark$
2. $E_0 = CB_0 (32.4)$ (from Faraday's Law)
 $\frac{E_0}{\sqrt{2}} = \varkappa \frac{E_0}{\sqrt{12}}$
3. $B_0 = E_0 ... 0$ (E₀ (32.8) (from Ampere's Law)
Hhis is sansfied in Vacuum because $c = 1$ \checkmark Page 10
Yes, it sanisties Maxwell's equations $\sqrt{E_0 ...}$

10 points

(equation # in textbook)

Consider an LRC series circuit driven with an ac source $v(t) = V_0 \cos(\omega t)$. You may use without proof the impedance and phase of an LRC series circuit.

(a): 5 points

Calculate the total energy dissipated in the resistor over one cycle if the system is driven at its resonant frequency.

at resonance frequency,

$$W_{0} = \frac{1}{VLC} (31.32)$$

$$T = \frac{2\pi}{W} = 2\pi \sqrt{LC}$$

$$P = \frac{V^{2}}{W} (25.18)$$

$$Total energy = \int_{0}^{2\pi\sqrt{LC}} \frac{(V_{0}\cos(\frac{1}{\pi c}))^{2}}{R} dt$$

$$= \frac{V_{0}^{2}}{K} \int_{0}^{2\pi\sqrt{LC}} \cos^{2}(\frac{1}{\sqrt{LC}}) dt$$

$$= \frac{V_{0}^{2}}{K} \left[\frac{1}{2} + \sqrt{LC}\sin(\frac{2t}{\sqrt{LC}})\right]_{0}^{2\pi\sqrt{LC}}$$

$$= \frac{V_{0}^{2}}{K} \left(\frac{2\pi\sqrt{LC}}{Z}\right)$$

$$= \frac{V_{0}^{2}\pi\sqrt{LC}}{R}$$

(b): 5 points

Calculate the total energy dissipated in the resistor over one cycle if the system is driven at *twice* its resonant frequency. Is it greater than or less than the energy found in part (a)?

in this case,

$$\begin{split} & \psi = 2\psi_{0} \\ &= \frac{2}{\sqrt{LC}} \\ T = \frac{2\pi}{\omega} = \pi\sqrt{LC} \\ total energy = \int_{0}^{\pi\sqrt{LC}} \frac{(V_{0}\cos\left(\frac{2\pi}{LC}\right))^{2}}{R} dt \\ &= \frac{V_{0}^{2}}{R} \int_{0}^{\pi\sqrt{LC}} \cos^{2}\left(\frac{2\pi}{LC}\right) dt \\ &= \frac{V_{0}^{2}}{R} \left[\frac{1}{2} + \sqrt{LC}sin\left(\frac{4\pi}{LC}\right)\right] \Big|_{0}^{\pi\sqrt{LC}} \\ &= \frac{V_{0}^{2}}{R} \left(\frac{\pi\sqrt{LC}}{2}\right) \\ &= \frac{V_{0}^{2}\pi\sqrt{LC}}{2R} \end{split}$$

it is 1/2 of the energy found in part a., so it is less than the energy found in part a.