

- 1) Two thin, parallel, conducting sheets of dimension $D \times W$ are separated by a distance d as shown ($D \gg W \gg d$). The sheets each carry a linear current density K , one into the plane of the page, one out, as shown.

- 1a) (15 points) Use Faraday's law to calculate the self-inductance of the arrangement.

The Faraday loops will actually run perpendicular to the lines of \vec{B} . Take $d\vec{s}$ parallel to \vec{B}

$$\oint_B d\vec{s} \cdot \vec{B} = \mu_0 K D d$$

$$\oint_E d\vec{s} = \mu_0 K W$$

$$\oint_B d\vec{s} = \frac{\mu_0 K D}{W} d$$

$$E_i = -\frac{d\Phi_B}{dt}$$

$$E_i = -\frac{\mu_0 dD}{W} \frac{dI}{dt}$$

$$L = \left| \frac{E_i}{\frac{dI}{dt}} \right| \Rightarrow L = \frac{\mu_0 dD}{W}$$

- 1b) (10 pts) Verify your answer to the first part using energy considerations.

$$U_B = \frac{1}{2} \mu_0 B^2$$

$$U_B = \frac{1}{2} \frac{\mu_0}{W^2} I^2$$

$$U = \int U_B dV$$

$$U = \frac{1}{2} \frac{\mu_0}{W^2} I^2 dDW$$

$$U = \frac{1}{2} \frac{\mu_0 dD}{W} I^2 \quad \text{Compare } U = \frac{1}{2} L I^2$$

$$L = \frac{\mu_0 dD}{W}$$

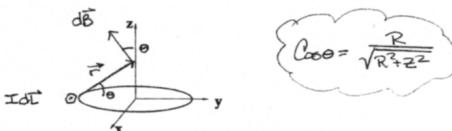
- 1c) (5 pts) Define \hat{L} as the inductance per unit length (measured along the current) and \hat{C} as the capacitance per unit length. Calculate $\sqrt{\hat{L}\hat{C}}$. This quantity plays an important role in the practical evaluation of transmission lines. Care to guess what it is?

$$\hat{L} = \frac{\mu_0 d}{W} \quad \hat{C} = \frac{\epsilon_0 W}{d}$$

$$\frac{1}{\sqrt{\hat{L}\hat{C}}} = \frac{1}{\sqrt{\mu_0 \epsilon_0}}$$

Essentially a fundamental constant.
With dimensionality $\frac{L}{c}$
... (the speed of light in Vacuum)

Symmetry:
 $\vec{B} = B_z \hat{z}$
 $B_z = \int d\vec{B}_z$
 $dB_z = dB \cos \theta$



$$B_z = \frac{R}{\sqrt{R^2 + z^2}}$$

- 2a) (10 points) A circular, conducting loop of radius R lies in the x,y -plane, centered on the origin. A current I flows through the loop such that at $z = +R$, the current is headed in the $+\hat{y}$ direction, and at $z = -R$ the current is headed in the $-\hat{y}$ direction. Derive the resultant magnetic field (magnitude and direction) at every point along the x -axis.

$$d\vec{B} = \frac{\mu_0}{4\pi} \frac{Id\vec{z} \times \hat{z}}{r^2}$$

$$dB = \frac{\mu_0 I}{4\pi} \frac{|d\vec{z}| \times \hat{z}}{r^2}$$

$$dB = \frac{\mu_0 I}{4\pi} \frac{dr}{(R^2 + z^2)^{3/2}}$$

$$dB_z = dB \cos \theta$$

$$dB_z = \frac{\mu_0 I R dr}{4\pi (R^2 + z^2)^{3/2}}$$

$$B_z = \frac{\mu_0 I R}{4\pi (R^2 + z^2)^{3/2}} \int dl$$

$$\vec{B} = \frac{\mu_0 I R^2}{2(R^2 + z^2)^{3/2}} \hat{z}$$

- 2c) (10 points) Now let's replace the thin ring of part b with a washer that extends from $r = a$ to $r = b$. Charge is distributed over the washer with a surface charge density

$$\sigma(r) = \frac{q ab}{2\pi(b-a)} \frac{1}{r^3}$$

and it rotates with a constant angular velocity ω . Find the magnitude and direction of the magnetic field produced at every point on the z -axis.

Superpose an infinite number of infinitesimal rings

$$d\vec{B} = \frac{\mu_0 \omega}{2} \frac{q ab}{2\pi(b-a)} \frac{dr}{(r^2 + z^2)^{3/2}}$$

$$d\vec{B} = \frac{\mu_0 \omega q ab}{4\pi(b-a)} \frac{z dz}{r^2 \sin^2 \theta} \frac{\cos^2 \theta}{z^3}$$

$$\int d\vec{B} = \frac{\mu_0 \omega q ab}{4\pi(b-a) z^2} \int_0^{a_0} dz \cos^2 \theta$$

$$\vec{B} = \frac{\mu_0 q ab \omega}{4\pi(b-a) z^2} \left. \sin \theta \right|_{a_0}^{a_0}$$

$$r = z \tan \theta$$

$$dr = \frac{z dz}{\cos^2 \theta}$$

$$\sqrt{r^2 + z^2} = \frac{z}{\cos \theta}$$

$$\vec{B} = \frac{\mu_0 q ab \omega}{4\pi(b-a) z^2} \left[\frac{b}{\sqrt{b^2 + z^2}} - \frac{a}{\sqrt{a^2 + z^2}} \right]$$

- 2b) (10 points) Let's replace the loop with an infinitesimally-thin, uniform ring of electric charge that extends from r to $r + dr$. If the surface charge-density on the ring is given by σ and the ring rotates about the z -axis with a constant angular velocity ω (recall, the direction of ω is obtained by using the right-hand rule with the physical motion of points on the ring) find the magnitude and direction of the (infinitesimal) magnetic field produced at every point along the z -axis.

Assume $\hat{\omega} = \hat{z} \Rightarrow$ then our ring is well-described by the work we did in part a

$$d\vec{B} = \frac{\mu_0 dI r^2}{2(r^2 + z^2)^{3/2}} \hat{z} \quad \leftarrow dI = \frac{dq}{l} = \frac{\sigma dA}{2\pi l} = \sigma \omega r dr$$

$$d\vec{B} = \frac{\mu_0 \sigma \omega r^2 dr}{2(r^2 + z^2)^{3/2}} \hat{z}$$

$$d\vec{B} = \frac{\mu_0 \omega}{2} \frac{\sigma r^3 dr}{(r^2 + z^2)^{3/2}}$$