		Danny Nguyer
. ,	Physics 1A - Winter 2016 Lecture 3	
	Final Exam	

Problem 1.

Shot = h

L = h/shot

Seeper sering hind

h

Shot B

Nond

h

Shot B

Nond

h

Shot B

Nond

h

Shot B

Nond

Two identical boxes slide down different ramps pictured above having started at rest from their tops. The boxes start at the same height. The coefficient of kinetic friction  $\mu_k$  between the boxes and the ramps is the same in both cases, and both ramps are fixed to the ground. The coefficient of static friction is not large enough to prevent the blocks from sliding down the ramps.

Let box A be the box on the left, and let box B be the box on the right. Consider the following statements:

- I. The speed of box A is greater than the speed of box B when they reach the bottoms of their ramps.
- II. The speed of box A is less than the speed of box B when they reach the bottoms of their ramps.
- III. The speed of box A is the same as the speed of box B when they reach the bottoms of their ramps.

## Questions.

- (a) Use force methods (energy methods not allowed!) to determine which if these statements is true when  $\mu_k \neq 0$ . You'll find it useful to name your own relevant variables so that you can answer this question mathematically.
- (b) Use energy methods to determine which if these statements is true when  $\mu_k \neq 0$ . You'll find it useful to name your own relevant variables so that you can answer this question mathematically.
- (c) Does the answer change when  $\mu_k = 0$ ? Justify using your answers from the previous parts.

a) Box A: 
$$Z = N - mg\cos\theta = 0$$
 $N = mg\cos\theta$ 
 $Z = mg\sin\theta - MkN = ma_A$ 
 $mg\sin\theta - mg\cos\theta \cdot Mkc = ma_A$ 
 $a_A = g\sin\theta - g\cos\theta \cdot Mkc$ 

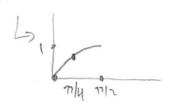
Box B:  $Z = N - mg\cos\theta = 0$ 
 $N = mg\cos\theta \cdot Mkc$ 
 $Z = mg\sin\theta - mg\cos\theta \cdot Mkc$ 
 $Z = mg\sin\theta - mg\cos\theta \cdot Mkc$ 
 $Z = mg\sin\theta - mg\cos\theta \cdot Mkc = ma_B$ 
 $mg\sin\theta - mg\cos\theta \cdot Mkc = ma_B$ 
 $mg\sin\theta - mg\cos\theta \cdot Mkc = ma_B$ 
 $mg\sin\theta - mg\cos\theta \cdot Mkc = ma_B$ 

Velocity A:  

$$V_{fA} = \left(2\left(g\sin\theta - g\cos\theta \cdot M\kappa\right) \cdot \frac{h}{\sin\theta}\right)^{1/2}$$
Velocity B:  $V_{fB}^2 = N_0^2 + 2ax$ 

$$V_{fB} = \left(2\left(g\sin\theta - g\cos\theta \cdot M\kappa\right) \cdot \frac{h}{\sin\theta}\right)^{1/2}$$

since block A is on a steeper incline, 0 > d. and cost And since sin &



less than I since the denominator (Sin & -cos & Mx) should be greater than

(sind-cosd Mic) because sind-cosomic > sind-cosd Mic due to the

B the fruction would be a small number/small number whereas for block A the fruction would be a small number/small number whereas for block A the fruction would be a shipply by number a number less than I,

So, velocity block A > velocity block 13,

b) block A

mgh = \frac{1}{2} mVA^2 + (MKN) \cdot \frac{h}{\sin \text{Sin } \text{ }}

gh = \frac{1}{2} VA^2 + MKg cos \text{ } \cdot \text{ } \frac{h}{\sin \text{ } \text{ }}

block B

 $mgh = \frac{1}{2} mV_B^2 + MICN \cdot \frac{h}{sm2}$   $gh = \frac{1}{2} V_B^2 + MICG \cos x$  Sind

Note that both blocks start with the same potential energy but they block b has a smaller normal fine than block A because cos & ccosa. Also, notice that block b traveres through a larger path, because sind & sind Sind >  $\frac{h}{\sin d}$  >  $\frac{h}{\sin d}$  >  $\frac{h}{\sin d}$  >  $\frac{h}{\sin d}$  > therefore, since block B has a larger normal fine and a longer path, the work done by father on block B> work done by father on block B> work done by father to block B> work done by father to translational.

Su, UB = Va.

In part A we can ignore the Arithmal force and obtain  $\alpha_A = g s in \theta$ ,  $\alpha_B = g s in \theta$ .

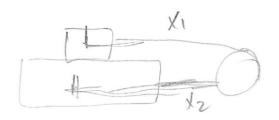
So, Velocities are clearly the same.

in part B we see that we can ignore the work done by fastown and obtain

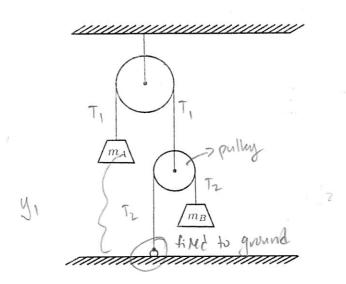
$$mgh = \frac{1}{2}mV_A^2$$

$$V_A = \sqrt{2g_A}$$

So velocities are the same.



Problem 2.



Consider the apparatus above. All pulleys and ropes are massless.

- (a) In limit  $m_B \to 0$ , would you expect the magnitude of the acceleration of mass A to be greater than, equal to, or less than g? Explain using physical reasoning.
- (b) Determine an expression for the acceleration of mass A in terms of the given variables.
- (c) Determine an expression for the tension in the top rope in terms of the given variables.
- (d) Does your mathematical answer in part (b) agree with your answer in part (a)? Explicitly verify this mathematically. If the answers don't agree, you should consider re-evaluating either your math, or your intuition, or both.
- a) I would expect amoss A = g. Since  $M_B \to O$ , the entire right size of the system will have a total of 0 mass. So, the tension in the rope connected to A would only be connected to mass A, meaning the tension M the rope will be O, so mass A will full with accelerating
- b) constraint:  $A_A = -A_{pully}$ ,  $T_1 = 2T_2$ mass  $B \sim T_2 m_B g = m_B a_B$ mass  $A \sim T_1 m_A g = m_A A_A$   $T_1 = 2T_2$   $T_2 = m_B g = m_B a_B$   $T_3 = m_B a_B$   $T_4 = m_A A_A$

Also, note that because of the bottom pully has a mass attached to the grand, length of rope = y1+nv +y1-ymb ~> Amb = 2. Apully

Equations: 
$$A_9 = -A_{pulley}$$
 $A_{13} = 2A_{pulley}$ 
 $T_1 = 2T_2$ 
 $T_2 - m_{Bg} = m_{I3}A_{I3}$ 
 $T_1 - m_{Ag} = m_{A}A_{A}$ 

$$T_{2} = m_{B}A_{B} + m_{B}g$$

$$A_{B} = \lambda A_{pulley} = -\lambda A_{a}$$

$$T_{2} = -\lambda m_{B}A_{A} + m_{B}g$$

$$L_{>} 2(-\lambda m_{B}A_{a} + m_{B}g) - m_{A}g = m_{A}A_{A}$$

$$\lambda m_{B}g - m_{A}g = A_{A}(\mu m_{B} + m_{A})$$

$$A_{A} = \frac{2m_{B}g - m_{A}g}{\mu m_{B} + m_{A}}$$

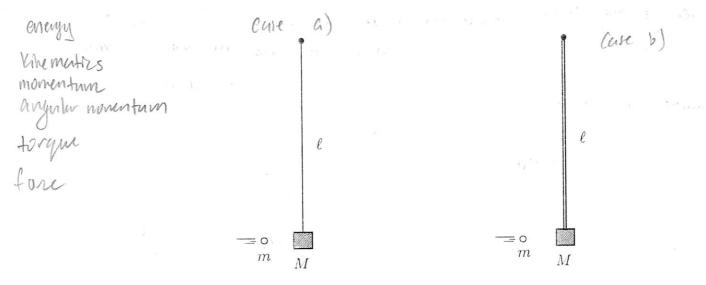
C) 
$$T_1 = m_A A_A + m_A g$$

$$T_1 = m_A \left( \frac{2 m_B g - m_A g}{4 m_B + m_A} \right) + m_A g$$

As 
$$M_B \rightarrow 0$$
,  $A_A = \frac{2(0) - M_A g}{4(0) + M_A g} = -\frac{M_A g}{M_A} = -\frac{g}{0 + M_A g}$ .

Also, from part C, as  $M_B \rightarrow 0$ ,  $T_1 = M_A \left( \frac{0 - M_A g}{0 + M_A} \right) + M_A g$ 

#### Problem 3.



In the diagram on the left, a small block of mass M is connected to a massless string of length  $\ell$  that is frictionlessly pivoted at its end so it can spin around in a vertical circle. Which be In the diagram on the right, an identical block of mass M is connected to a rigid, massless rod of length  $\ell$  that is frictionlessly pivoted at its end so it can spin around in a vertical circle. In both cases, the block starts out hanging at rest, and a clay pellet of mass m is fired horizontally at the block and gets lodged inside.

- (a) In which case would you expect the pellet needs to be shot with a higher speed for the block to move all the way around in a vertical circle with radius  $\ell$ ?
- (b) In the case on the left, with what speed does the pellet need to be shot at the block so the block will make it all the way around in a vertical circle of radius  $\ell$ ?
- (c) In the case on the right, with what speed does the pellet need to be shot at the block so the block will make it all the way around in a vertical circle of radius  $\ell$ ?
- (d) According to your answers to parts (b) and (c), in which case does the speed of the pellet need to be greater? Does your mathematics agree with your intuition from part (a)?
- a) I believe case A will need to have a pollet shot a higher speed. For case 13, the pellet and block system needs to Simply just past the maximum vertical height; after reaching this point gravity will simply let the system full down, they completing the vertical circle. In care (A) however, the system must maintain a constant contripetal acceleration cet  $\frac{V^2}{V}$  to complete the anle.

b) Because the collision between the bullet and block is extremely short, we can say that momentum is conserved she the impulse, which is equal to the change of momentum, is a because  $J = F \cdot time$  where time is close to 0,  $\sim > \frac{dP}{dt} = 2F \cdot dt \approx 0$ 

After collision, the pair is traveling at velocity Uf. At peak:

$$\frac{|\langle \mathcal{E}_{i} \rangle - |\mathcal{P}_{\xi} + |\mathcal{K}_{\xi}|}{\frac{1}{2} (m+m) V_{\xi}^{2}} = \frac{1}{2} (m+m) V_{\xi}^{2} + \frac{1}{2} (m+m) V_{\xi}^{2}$$

$$\frac{1}{2} V_{\xi}^{2} = \frac{1}{2} V_{\xi}^{2} + \frac{1}{2} V_{\xi}^{2}$$

$$\frac{1}{2} V_{\xi}^{2} = \frac{1}{2} V_{\xi}^{2} + \frac{1}{2} V_{\xi}^{2}$$

At top of circle.

mg is providing the force regular to keep System morty

=) my = 
$$mac = m \cdot v^2$$

$$\frac{1}{2}Vf^{2} = 2gL + \frac{1}{2} \cdot (\overline{\Gamma Lg})^{2}$$

$$Vf^{2} = UgL + Lg$$

$$Vf = \overline{\Gamma J Lg}$$

The bullet must be shot with a velocity of

c) In the case on the right, the bullet-block system shiply has to reach the top prak of the carle.

So, usty the same logic in part (B) for momentum conservative,  $mVi = (m+M)V_f$  where  $V_f$  is the velocity that the pair is moving with,

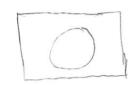
To reach the top and nothing else,

$$\frac{1}{2} (m+m)V_f^2 = (m+m)gH$$

$$V_f = \sqrt{2}g \cdot 2L$$

$$V_f = \frac{(m+m)}{m} V_f = \frac{(m+m) \cdot 2\sqrt{g}}{m}$$

d) From the answers obtained in parts B and C, the velocity for the non-nyill string case needs to be higher. This result does agree with my intrition since  $\frac{(m+m)}{m}$ .  $\overline{15gL} > \frac{(m+m)}{m}$ .  $2\overline{1gL}$ .



## Problem 4.

A large, uniform solid disk of mass M and radius R initially spins on the surface of a flat, frictionless surface at an angular speed  $\omega$ . Its center of mass is initially at rest relative to the table. Recall that the moment of inertia of a uniform solid disk for rotations about an axis passing through its center of mass and perpendicular to its face is  $MR^2/2$ .

Josh and Nancy are initially standing diametrically opposite one another on the edge of the disk (so the initial distance between them is 2R). Next they walk directly toward one another along the diameter joining them until they meet at the disk's center.

They both move with the same speed as a function of time and they both have mass (c/4)M where c is a unitless constant. Let P denote the total momentum of the Josh + Nancy + disk system. Let  $L_{\parallel}$  denote the angular momentum of the Josh +Nancy + disk system in the direction parallel to the axis of rotation.

- (a) Is P conserved as Josh and Nancy walk to the center?
- (b) What is the motion of the center of mass of the disk as Josh and Nancy walk to the center?
- (c) Is  $L_{\parallel}$  conserved as Josh and Nancy walk to the center?
- (d) Determine the angular speed of the disk when Josh and Nancy are at its center.
- (e) Is the mechanical energy of the system conserved as Josh and Nancy walk to the center? If it is conserved, prove it. If not, compute the change in mechanical energy and show that it's nonzero. In both cases, give physical reasoning to explain why your answer makes sense as well. If it doesn't make sense, you may consider re-evaluating either your intuition about this scenario, or your math, or both.
- (f) What would you expect the answer to parts (d) and (e) would be in the limit  $c \to 0$ ? Do your mathematical answers agree with these expectations?

a)

Josh Namy

Josh Namy

disk my

There are no net external forces acting an Josh, Nany, and the disk, so their momentum is conserved. Also, she all forces are in the Z-direction, and  $T = V \times F$ , we know that there are no net torque in the Z-direction but there will be in the other directions. So Lz is conserved but Lx and Ly may not be. So total momentum is not conserved.

b) since the disk was unitarally solid, the counter of mass was initially at the centre, then because the two people walk with the same speed and have the same mass, and began equiditant from the centre, the centre of mass does not

change as the pair moves toward the center, Also, Fertinet = 0 = M. acm. Since acm = 0, and the disk is at vest relative to table, Vcm = 0 => position of center of mass is constant.

External torques in the Z-direction in part A. There are no net suffered torques in the Z-direction became all forces act on the Z-direction. So since I = VXF, VXF is not in Z-direction, and thus

IZ = 0 => L2 = L1, is conserved because \(\frac{2}{ct}\) \(\frac{2}{ct}\).

d) initial anymer speed > w initial moment of herrin = \frac{1}{2}MR^2 + 2(\frac{1}{4}M)R^2

And moment of Metria = { MR2 since both people are at the center

su, L11 conserved

$$= \sum_{n=1}^{\infty} \left( \frac{1}{2} MR^2 + \frac{c}{2} MR^2 + \frac{c}{2} MR^2 \right) W = \frac{1}{2} MR^2 \cdot W_{final}$$

$$= \frac{1}{2} MR^2 + \frac{c}{2} MR^2$$

$$= \frac{1}{2} MR^2$$

$$= \frac{1}{2} MR^2$$

$$= \frac{1}{2} MR^2$$

$$= \frac{1}{2} MR^2$$

- e) Mechanical energy 15 conserved. As the two people are walking towards the center, they exect equal but opposite forces onto the disk. But because they are exerting equal but opposite forces, the net work done is a. Moreover, the other forces acting on this system are growing and normal forces. But the normal force is preparatively to motion, so work due to normal forces =0. As such, since the work of noncommative forces is a, mediantial 15 energy must be conserved.
- f) As C>0, the two people become massless, so angular velocity should remain unchanged, and mechanical energy should be the same as well. In fact our answers support this.

In part (D) No WAMAI = (1+c) WO WAMAI = WO.

In part (e) ~> as c>0
mechanizal energy is still conserved.

### Problem 5.

- (a) Consider a projectile launched at speed v angle  $\theta$  relative to the horizontal on flat ground near the surface of a planet with gravitational acceleration  $g_P$ . Derive an expression for the range  $r_P$  of the projectile on this planet.
- (b) If an object of mass m is a distance r away from the center of a planet of mass M, then it experiences an attractive gravitational force of magnitude

$$F = \frac{GMm}{r^2} \tag{1}$$

where G is Newton's gravitational constant. Using Newton's Second Law to set this equal to the object's mass times the magnitude a of its acceleration, we find that the object's gravitational acceleration is independent of its mass, but depends on G, M and r:

$$a = \frac{GM}{r^2} \tag{2}$$

In other words, it depends only on a fundamental physical constant, the mass of the planet, and the distance to the planet's center. If the object is at a height h above the planet's surface, and if the planet's radius is R, then the gravitational acceleration becomes

$$a = \frac{GM}{(R+h)^2} = \frac{GM}{R^2} \left(1 + \frac{h}{R}\right)^{-2} \tag{3}$$
 What is the Taylor expansion of the acceleration due to gravity in the variable  $x = h/R$ 

about x = 0 including only the first three nonzero terms?

Useful observations. You should find that the first non-vanishing order equals  $GM/R^2$ . When you are close to the surface of the planet, only this first nonzero term is significant because x = h/R will be extremely small, so this expression gives the acceleration due to gravity near the planet's surface.

$$g_P = \frac{GM}{R^2}. (4)$$

In the case of Earth, one can for example show that by using this term to compute the acceleration due to gravity, the quantity  $GM_E/R_E^2$  gives a value very close to g, where  $M_E$  is the Earth's mass, and  $R_E$  is the Earth's radius. In other words, one can predict the acceleration due to gravity near the Earth's surface using its mass and its radius!

(c) Alice is on planet A whose mass is  $M_E/2$  and whose radius is  $R_E$ . Bob is on planet B whose mass is  $\sqrt{3}M_E$  and whose radius is  $\sqrt{2}R_E$ .

Alice throws a ball at a speed v at a certain angle relative to the ground that maximizes the range of the thrown object. Bob throws a ball at speed  $\sqrt{2}v$  and at an angle  $\theta$  relative to the ground, and the ball ends up having the same range as Alice's ball.

At what angle  $\theta$  did Bob throw the ball?

Very construction of the state of the projective to reach the ground again, and we know that the distance to the 
$$\frac{1}{2}$$

and we know that the distance traveled = 
$$X = Vox \cdot t$$
,

range =  $Vcos \theta \cdot 2Vsih \theta = 2V^2 cos \theta sih \theta$ 
 $g_P$ 

b) Taylor Expansion
$$f(x) = f(a) + \frac{f'(a)(x-a)}{1!} + \frac{f''(a)(x-a)}{2!} + \dots$$

$$agravitational = \frac{GM}{(12+h)^2} = \frac{GM}{R^2} \left(1 + \frac{h}{r}\right)^{-2}$$

$$a' = \frac{-2 GM}{R^2} \left(1 + \frac{h}{r}\right)^{-3}$$

$$a' = \frac{-2 GM}{R^2} \left(1 + \frac{h}{r}\right)^{-3}$$

$$a'' = \frac{66M}{R^2} \left( 1 + \frac{h}{r} \right)^{-4}$$

Taylor Expansion about 
$$x=0$$
 yields 
$$\frac{GM}{R^2} - \frac{2GM}{R^2} \left( x-0 \right) + \frac{GGM}{R^2} \left( \frac{x-0}{2} \right)^2 = \left[ \frac{GM}{R^2} \left( 1 - 2x + 3x^2 \right) \right]$$

plane+ B- mass 13 Me radius TZ RE Extra Space

gravity on plant 
$$A = \frac{6M\epsilon}{\frac{7}{2}} = \frac{1}{2}9$$

gravity in earth =  $g = \frac{GM\epsilon}{R\epsilon^2}$ 

Alice:  $\sqrt{6}$  range =  $\frac{2v^2 \sin \theta \cos \theta}{\frac{1}{2}q} = \frac{4v^2 \sin \theta \cos \theta}{q} = \frac{2v^2 \sin \theta \cos \theta}{q}$ 

 $range = \frac{4V^2 sinkcosd}{\sqrt{3}q} = \frac{8V^2 sindcosd}{\sqrt{3}q} = \frac{4V^2 sin(2d)}{\sqrt{3}q}$ 

$$\frac{4v^{2}\sin(2d)}{\sqrt{3}g} = \frac{2v^{2}\sin(2G)}{g}$$

$$\sin(2d) = \frac{\sqrt{3}}{2}\sin(2G)$$

To maximize range, however, Alice must have thrown the ball at  $\Theta = 714$ 

$$X = d\sqrt{\cos\theta} \sin\theta$$

$$X' = d\sqrt{2} \cos\theta \sin\theta$$

$$X' = d\sqrt{2} \left(-\sin^2\theta + \cos^2\theta\right) = 0$$

$$X' = 3\sin^2\theta = \cos^2\theta \Rightarrow \theta = 7/4 \text{ to maximize}$$

varye

So, 
$$\sin(2d) = \frac{13}{2} \sinh(\frac{\pi}{2})$$
  
 $\sin(2d) = \frac{13}{2} \sinh(\frac{\pi}{2})$   
 $\sin(2d) = \frac{13}{2} \implies 2d = \frac{\pi}{3} \implies \alpha = \frac{\pi}{6}$