ECE170A Midterm Exam

Date: 11/6/18 4-5:50 pm

1.

Snell's law and lateral beam displacement What is the displacement of a laser beam passing through a glass plate of thickness 2 mm and refractive index 1.570 if the angle of incidence is 40°? (See Figure 1.14)

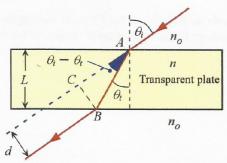


Figure 1.14 Lateral displacement of light passing obliquely through a transparent plate

2. Note: No worry about the spectroscopic ellipsometry measurements, the given photon energy is 1.5 eV in this problem

Complex refractive index Spectroscopic ellipsometry measurements on a germanium crystal at a photon energy of 1.5 eV show that the real and imaginary parts of the complex relative permittivity are 21.56 and 2.772 respectively. Find the complex refractive index. What is the reflectance and absorption coefficient at this wavelength? How do your calculations match with the experimental values of n = 4.653 and K = 0.298, R = 0.419 and $\alpha = 4.53 \times 10^6$ m⁻¹?

3.

Fabry-Perot optical cavity from a ruby crystal Consider a ruby crystal of diameter 1 cm and length 10 cm. The refractive index is 1.78. The ends have been silvered and the reflectances are 0.99 and 0.95 each. What is the nearest mode number that corresponds to a radiation of wavelength 694.3 nm? What is the actual wavelength of the mode closest to 694.3 nm? What is the mode separation in frequency and wavelength? What are the finesse F and Q-factors for the cavity?

4.

A multimode fiber Consider a multimode fiber with a core diameter of 100 μ m, core refractive index of 1.4750, and a cladding refractive index of 1.4550 both at 850 nm. Consider operating this fiber at λ = 850 nm. (a) Calculate the *V*-number for the fiber and estimate the number of modes. (b) Calculate the wavelength beyond which the fiber becomes single mode. (c) Calculate the numerical aperture. (d) Calculate the maximum acceptance angle. (e) Calculate the modal dispersion $\Delta \tau$ and hence the bit rate × distance product.

Dispersion at zero dispersion coefficient Since the spread in the group delay $\Delta \tau$ along a fiber depends on the $\Delta \lambda$, the linewidth of the source, we can expand $\Delta \tau$ as a Taylor series in $\Delta \lambda$. Consider the expansion at $\lambda = \lambda_0$ where $D_{ch} = 0$. The first term with $\Delta \lambda$ would have $d\Delta \tau/d\lambda$ as a coefficient, that is D_{ch} , and at λ_0 this will be zero; but not the second term with $(\Delta \lambda)^2$ that has a differential, $d^2\Delta \tau/d\lambda^2$ or $dD_{ch}/d\lambda$. Thus, the dispersion at λ_0 would be controlled by the slope S_0 of D_{ch} vs. λ curve at λ_0 . Show that the chromatic dispersion at λ_0 is

$$\Delta \tau = \frac{L}{2} S_0 (\Delta \lambda)^2$$

A single mode fiber has a zero-dispersion at $\lambda_0 = 1310$ nm, dispersion slope $S_0 = 0.090$ ps nm² km. What is the dispersion for a laser with $\Delta \lambda = 1.5$ nm? What would control the dispersion?

1.	L=2mm, n=1.970, 0;=40° nd=1
	1. / RMM 2/
	d=Lomo: (1- Tomo?-sin20;)
	= 2mm siny 0° (1 - (05 40°) /1,5702 - snin 240)
	d=0.598 mm
7	200 C 3 T 3 T 3 T 4 T 1 3 C 1 C 2 C 2 C 2 C 2 C 2 C 2 C 2 C 2 C 2
1.	E=1.5eV, &=21.56-12.772, &=21.56, &=2.772
	$N = \sqrt{\frac{1}{2}} \left(\xi_{r}^{1} + \sqrt{\xi_{r}^{12} + \xi_{r}^{112}} \right)$
	11 - 1 2 - 1 - 1 - 1 - 1 - 1 - 1 - 1 - 1
	$=\sqrt{\frac{1}{2}(21.56+\sqrt{(21.56)^2+(2.772)^2}}$
	(n=4,65282)
	$Y = \int \frac{(\xi_{\gamma}^{1/2} + \xi_{\gamma}^{1/2})^{1/2} - \xi_{\gamma}^{1/2}}{2} = \int \frac{(21.56^{2} + 2.772^{2})^{1/2} - 21.56}{2}$
	N=n-iK
	(K=8.29788) N=4.65282-j0.29788
	$(N-1)^2+K^2$
	$R = \frac{(N-1)^2 + K^2}{(N+1)^2 + K^2}$
	$-(4.65282+1)^2+(0.29788)^2$
	$= \frac{(4.65282+1)^2 + (6.29788)^2}{(4.65282+1)^2 + (6.29788)^2}$
	(R=191918)
	(A) = 11337 am (Section 2)
	a=2kgK E=hv=> v=3.622.10"
	$\lambda = \frac{C}{2}$
	$\frac{2\pi}{k} = \frac{\Gamma}{V} = \frac{2\pi V}{C}$
	C1878188888=D1C=7586098/m1
	α=2h, K=2.7586098/m·0.29788
	a = 4.52.106 mil
	Desce a revenue la
	These are very close to the experimental values.

3.	d=lom, L=10cm, n=1,78, R=0,99, R=0,95, closest mode	
	to x= 694,3 nm?	
- Dr	$m = \frac{2nL}{\lambda} = \frac{2 \cdot 1.78 \cdot 16 cm}{694.3nm}$ $m = 512.746.65$	
	$M = \frac{2NL}{\lambda} = \frac{2NL}{694.3nm}$	
	=>(m=612747 is closed to 694,3nm	
	2nl 2/178/18	
	$\lambda_{m} = \frac{2nL}{m} = \frac{2\cdot1.78\cdot10_{cm}}{512747}$	
	-/ nm = 091, 2995 nm)	
	at m= 512747.	
	(w £ 2323)	
	Average reflectance is 2 = 0,97	
	and the same of th	
	Finesse F= 1-12 = 1-9/97 = (F=183.137)	
	NAC OCAL CALANAMAN AND AND AND AND AND AND AND AND AND A	
	Mode separentran DVm = 2n2 = 3.10 sm/s	
and the second s	· · · · · · · · · · · · · · · · · · ·	
	(DVm=8,427×108 +12)	
	100/2/01/2019	
	$\Delta \lambda = \frac{\lambda m^2}{2nL} - \frac{(694.2995 \text{ nm})^2}{2.1.78 \cdot 1000}$	
	- 5 H2 - 12 - 12 - 12 - 12 - 12 - 12 - 12 -	
	(DX = 1.354 pm (preometers))	
	Marks = 10	
	Q=m ₀ F	
	mo=512747	
	Q=812747.103.137 => (Q=5288/3187.34)	
	ACTIVE TO THE PROPERTY OF THE	
	(1-4351/20m2)	
	These aim very class to the expression with values	

M. a = 50 pm, neare = 1,4750, nelan = 1,4550 at \ = 850 nm a) $V = \frac{2\pi a}{\lambda} \sqrt{n_1^2 - n_2^2} = \frac{2\pi \cdot 50 \mu m}{850 nm} \cdot \sqrt{1.475^2 - 1.455^2}$ V=891.47 # of modes M = 1/2 = 89,472 M=4002,483 = (4002 modes) b) (utoft wavelength he when V= 2,409 = 200 /01,2-net $\lambda_{c} = \frac{2\pi\alpha}{2.405} \sqrt{n_{1}^{2} - n_{2}^{2}} = \frac{2\pi \cdot 50 \mu m}{2.405} \sqrt{1.475^{2} - 1.455^{2}}$ (hc = 31.62 pm) c) NA = \(\lambda_1^2 - \lambda_2^2 = \sqrt{1.475^2 - 1.455^2} \) NA = 0.24207 d) Maximum acceptance make dinax $8ma_{max} = \left(\frac{NA}{n_0}\right) \quad n_0 = 1$ amax = arcsintNA) = arcsin(0,24207) amox=19.01°) Total acceptance angle = 20max = 28.02° e) Modal dispersion 0 = 0,29 st ST=L. M1-M2 = L. 1975-1.455 DT=6,667.10-1/8/m.L BL = 0,28L - 0,28L - 0,28L - 0,29.6.667.10-11 2m BL= 1.293.100 500 m

 $\frac{\partial L}{L} = \left| \begin{array}{c} O(n + O_m) O \lambda \\ \approx \frac{1}{1!} \frac{o(OT)}{o(\lambda)} \left|_{\lambda = \lambda_0} O \lambda \right| + \frac{1}{2!} \frac{o^2(OT)}{o(\lambda)^2} \left|_{\lambda = \lambda_0} (O \lambda)^2 \right| Toylor exponeron \\ = O(n + O \lambda) + \frac{1}{2} \frac{o(OT)}{o(\lambda)} \left|_{\lambda = \lambda_0} (O \lambda)^2 \right| O(n = 0, so$ $\frac{\partial U}{L} = \frac{1}{2} \frac{\partial D_{eh}}{\partial \lambda^{2}} \Big|_{\lambda = \lambda_{0}} \cdot (\Delta \lambda)^{2} \qquad S_{0} = \frac{\partial D_{ch}}{\partial \lambda^{2}} \Big|_{\lambda = \lambda_{0}}$ DT = 250.(0)2 λ0=1310 nm, S0=0100 nm21hm, Δλ=1.5 nm DT = \frac{1}{2} \S_6 \cdot (DX)^2 = \frac{1}{2} \cdot (109 \frac{\rho_8}{\rm^2 \km} \cdot (1.5 \rm)^2 DT = 0110125 ps L= 1km => (DT = 0,10125 ps Dispersion would be controlled by sk and the Length of the fiber and the slope of Do at &= 1310 nm. Since sh=1,5mm, there is some wavelength vanadion in the light along the fiber, so some wordengths travel faster as they penetrate further into the cladding (whose n is lower) and so they travel faster, white others travel Slower, comming dispersion, due to shound So.